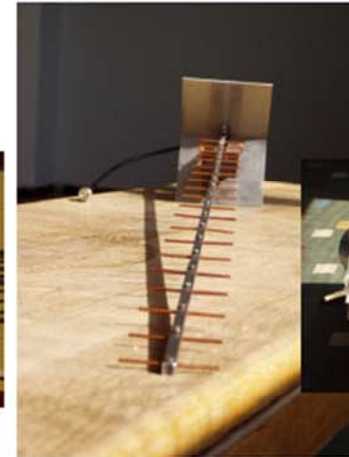


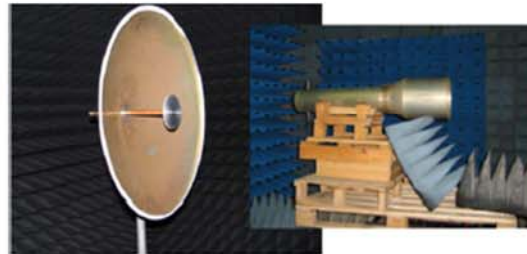
Antenna Basics 3

The “antenna family”

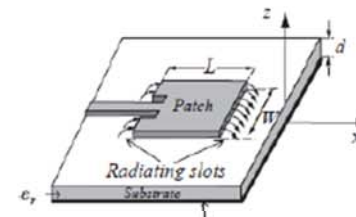
- **Wire antennas**
- Straight wire (dipole), loop, helix, ...



- **Aperture antennas**
- Horns, reflectors, lens



- **Planar (“microstrip” antennas)**
- Metallic patch on a grounded substrate

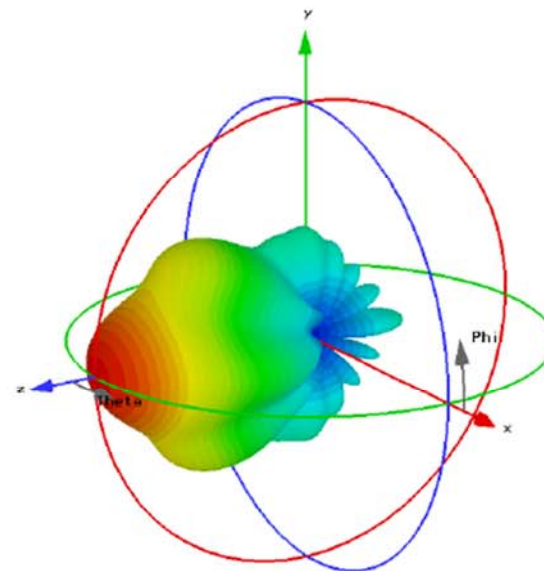
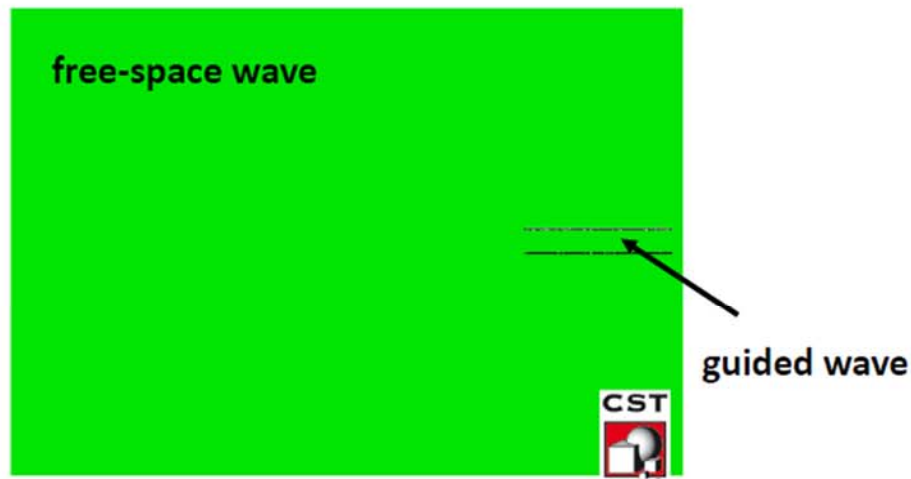


- **“Special” antennas**
- fractal geometry, on-chip antennas, mm-wave antennas, metamaterials, antennas for medical applications (implantable)...

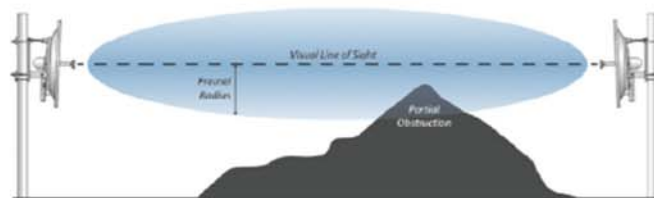


Antenna is...

- Passive filter, both in frequency and spatial domain
- Component transforming guided waves to waves in free space (that is actually 3D spherical waveguide)
- Very important component of the whole communication channel
- „The best amplifier“

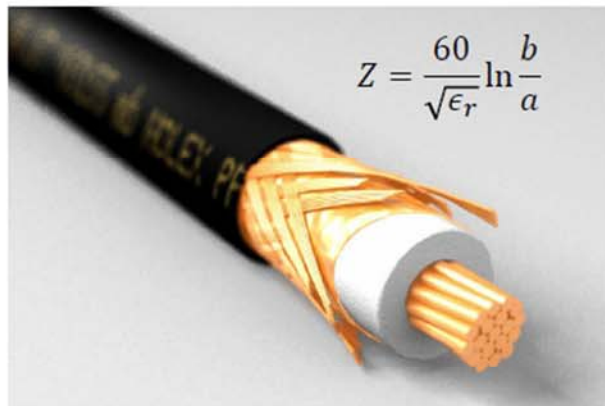


Antenna link – the low loss link



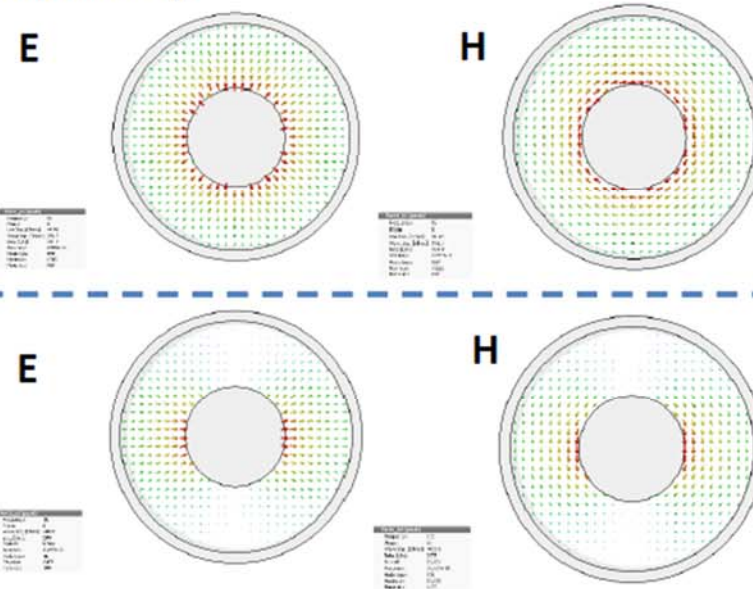
Waveguides

- With TEM wave (operation from DC): coaxial cable, multiconductor line
- Parameters not function of frequency! (no dispersion)



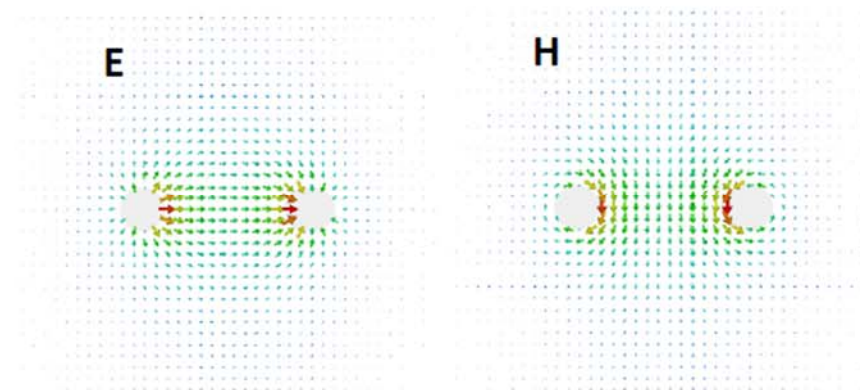
$$Z = \frac{60}{\sqrt{\epsilon_r}} \ln \frac{b}{a}$$

TEM



Quasi-TE11

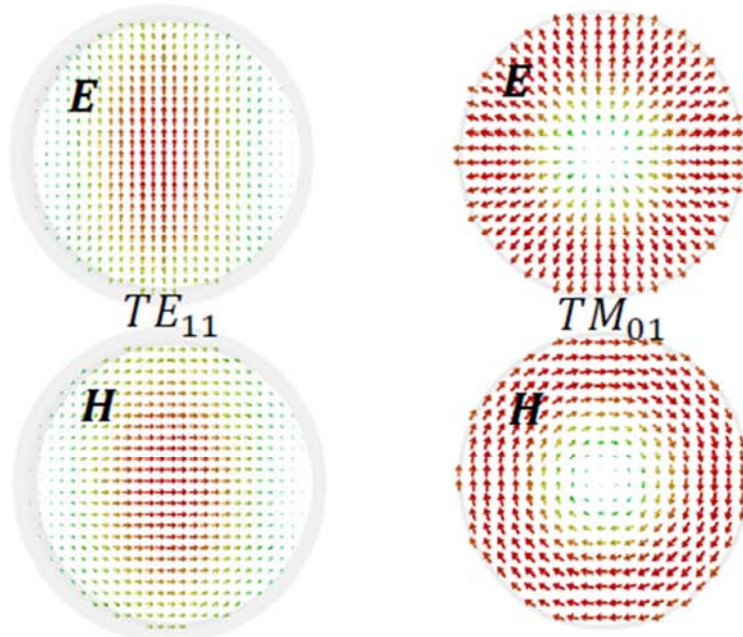
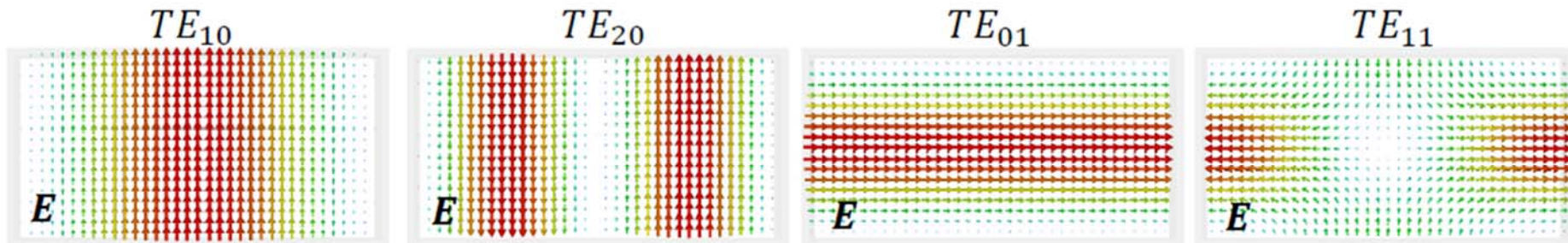
$$f_c \cong \frac{c_0}{\pi(a+b)\sqrt{\epsilon_r}}$$



$$Z = \frac{120}{\sqrt{\epsilon_r}} \cosh^{-1} \left(\frac{D}{d} \right)$$

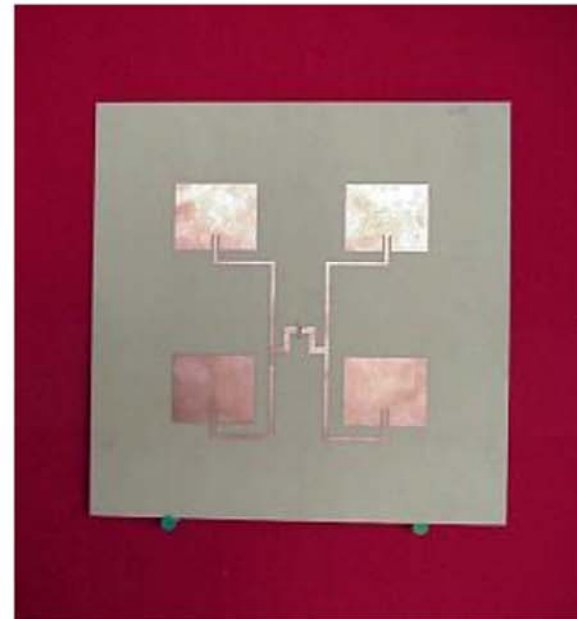
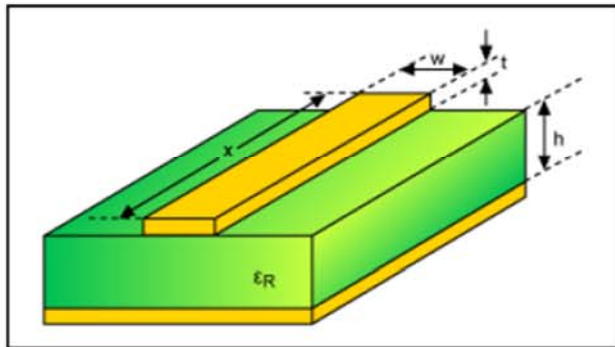
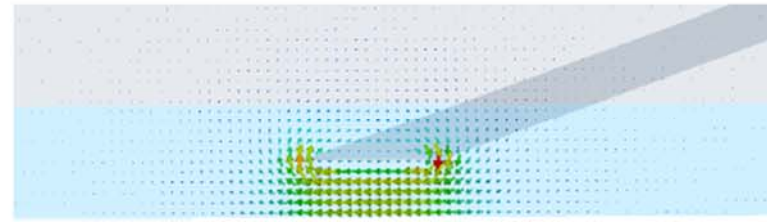
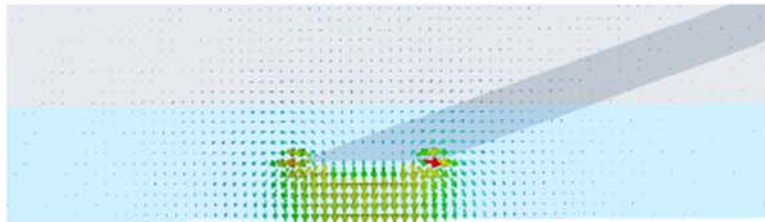
Waveguides - metallic

- With TE/TM modes (operation above cut-off frequency): metallic waveguides
- Parameters are function of frequency! (dispersion)

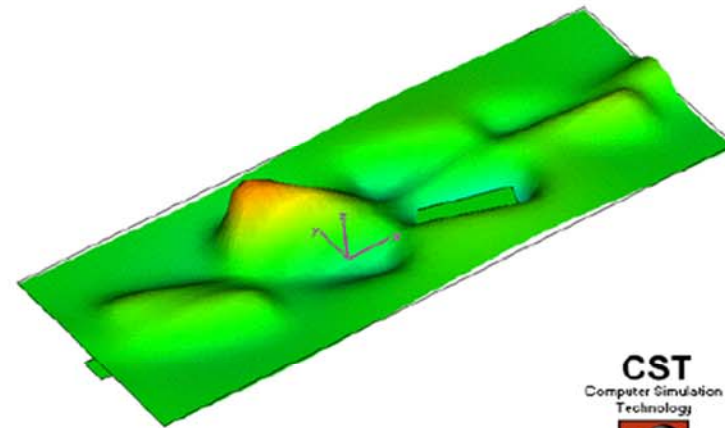
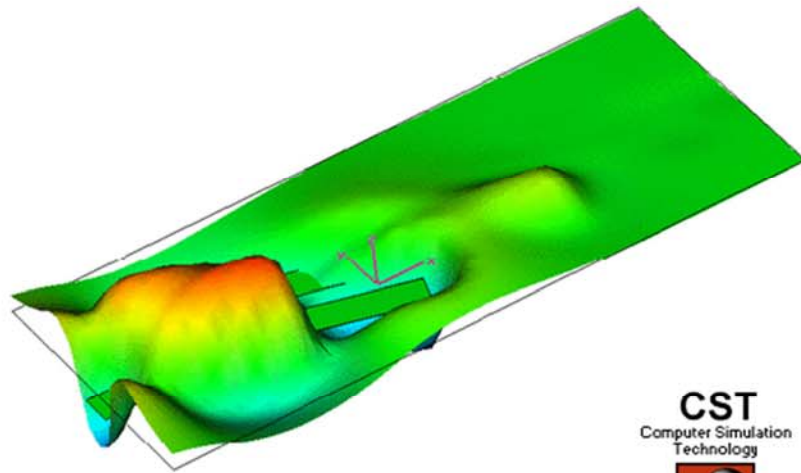
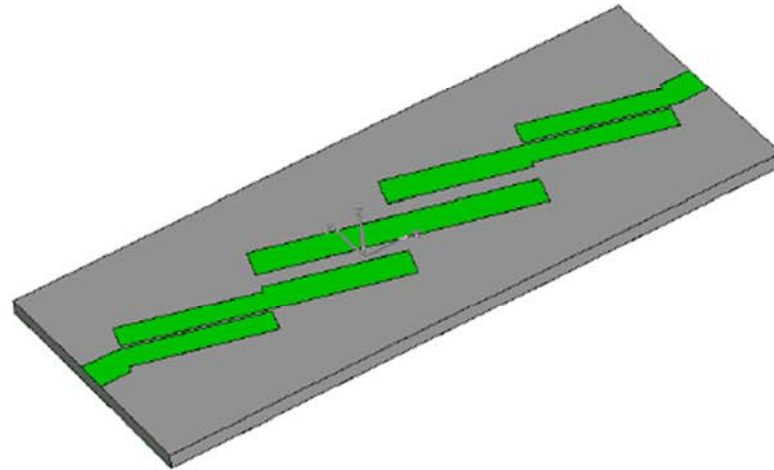


Waveguides - planar

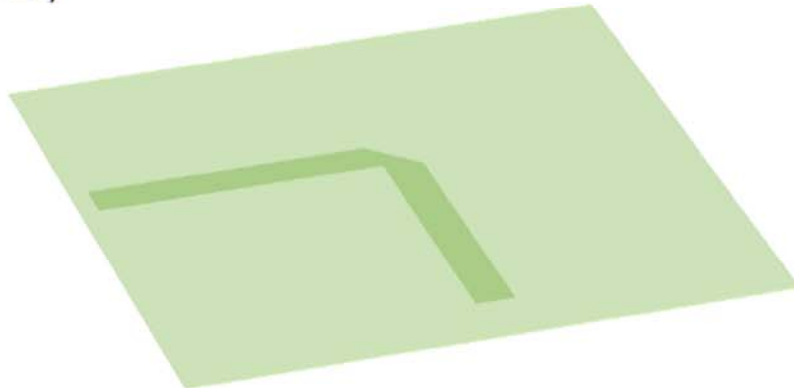
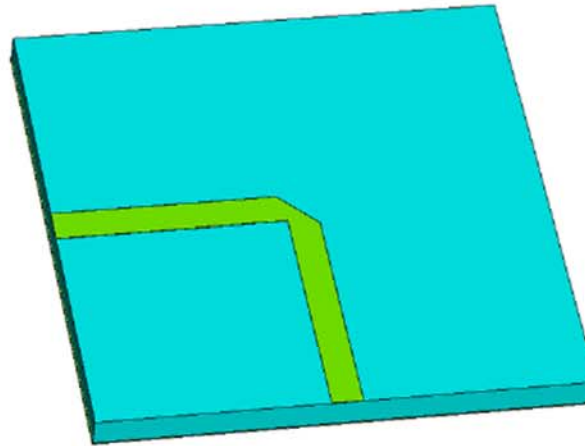
- With quasi-TEM wave (operation from DC possible): **microstrip**, stripline, ...
- Parameters are function of frequency! (slight dispersion)



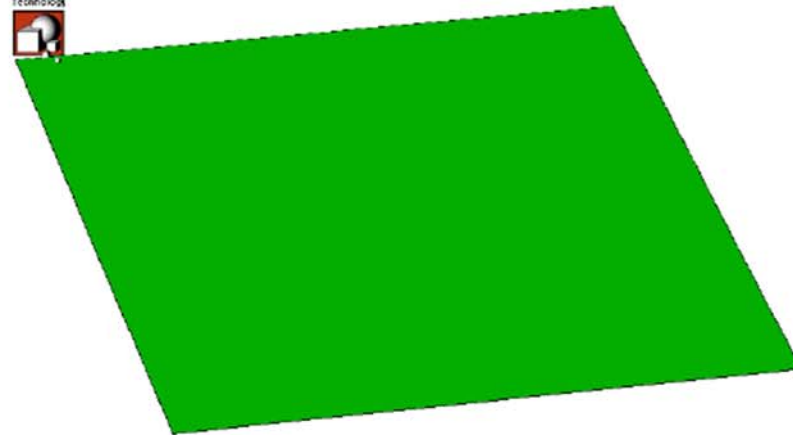
Waveguides – planar - microstrip



Waveguides – planar - microstrip



Type = Powerflow (peak)
Monitor = power (t=0..1(0.01)) [1]
Maximum = 213717 VA/m²
Sample = 1 / 100
Time = 0

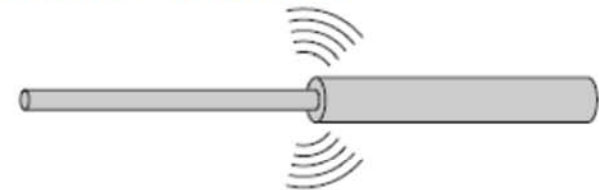
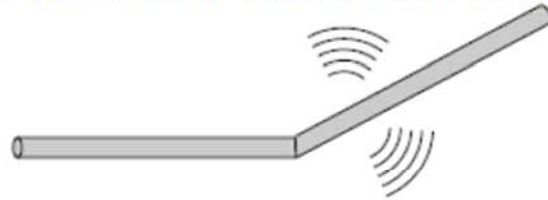


Type = E-Field (peak)
Monitor = e-field (t=0..1(0.01)) [1]
Component = z
Plane at z = 1.63389
Sample = 1 / 100
Time = 0

Radiation

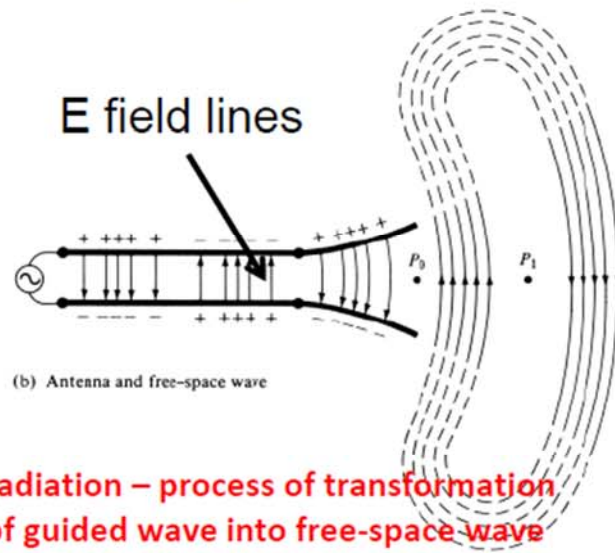
- Radiation is caused by *accelerating (deccelerating) charge – time varying current* (usually time-harmonic), discontinuities

$$\nabla \cdot \mathbf{J} = -\frac{\partial \rho}{\partial t}$$



- A stationary charge will not radiate EM waves – zero current → no magnetic field, no radiated power

$$P_{rad} = \frac{1}{2} \Re \oint \mathbf{E} \times \mathbf{H} d\mathbf{S}$$



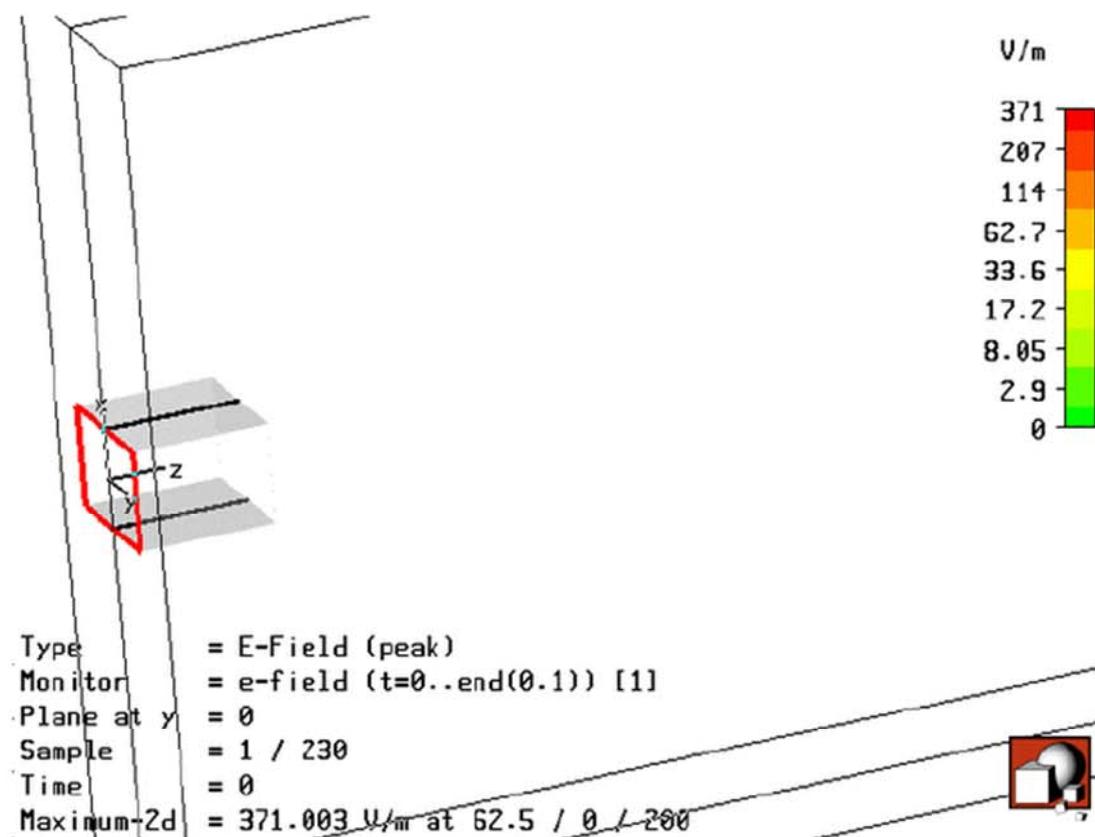
Electric field lines:

- Start on positive charges and end on negative charges
- Start on a positive charge and end at infinity
- Start at infinity and end on a negative charge
- Form closed loops neither starting or ending on any charge

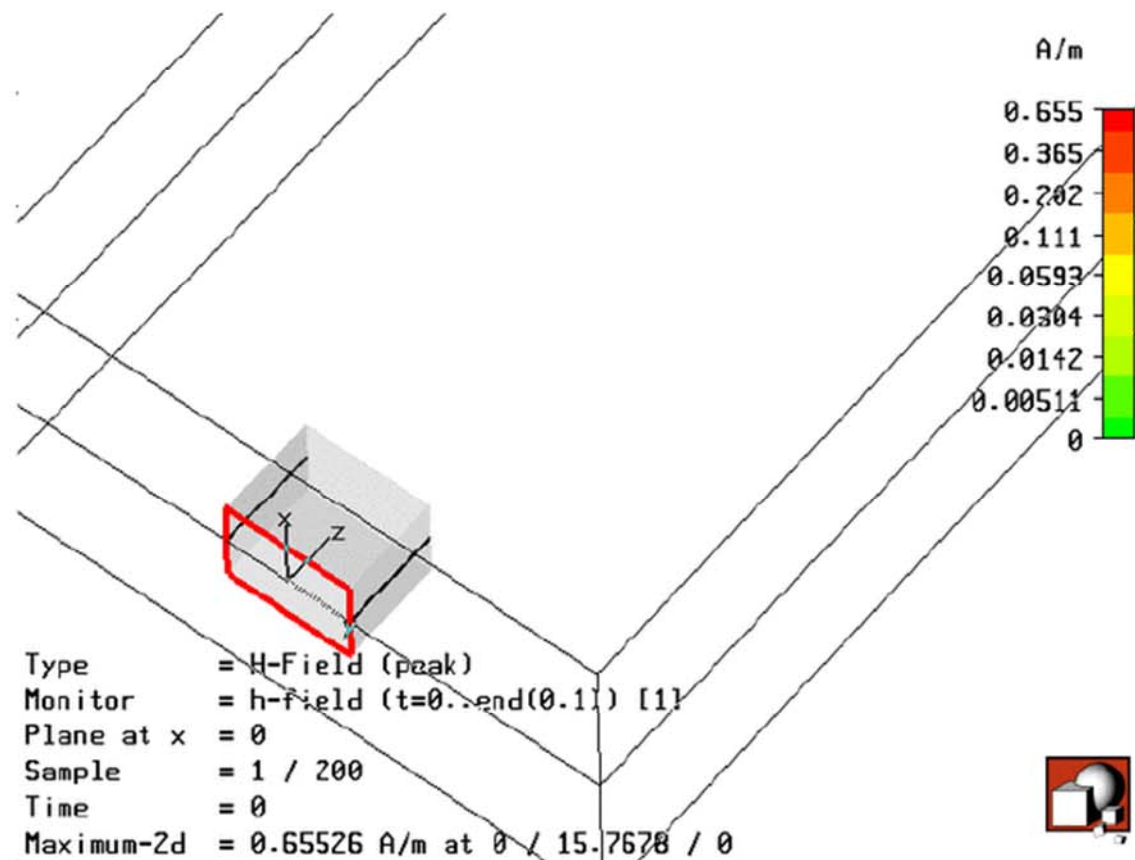
Magnetic field lines:

- Always form closed loops encircling current because physically there are no magnetic charges
- Mathematically it is often convenient to introduce equivalent magnetic charges and magnetic currents

Radiation - E

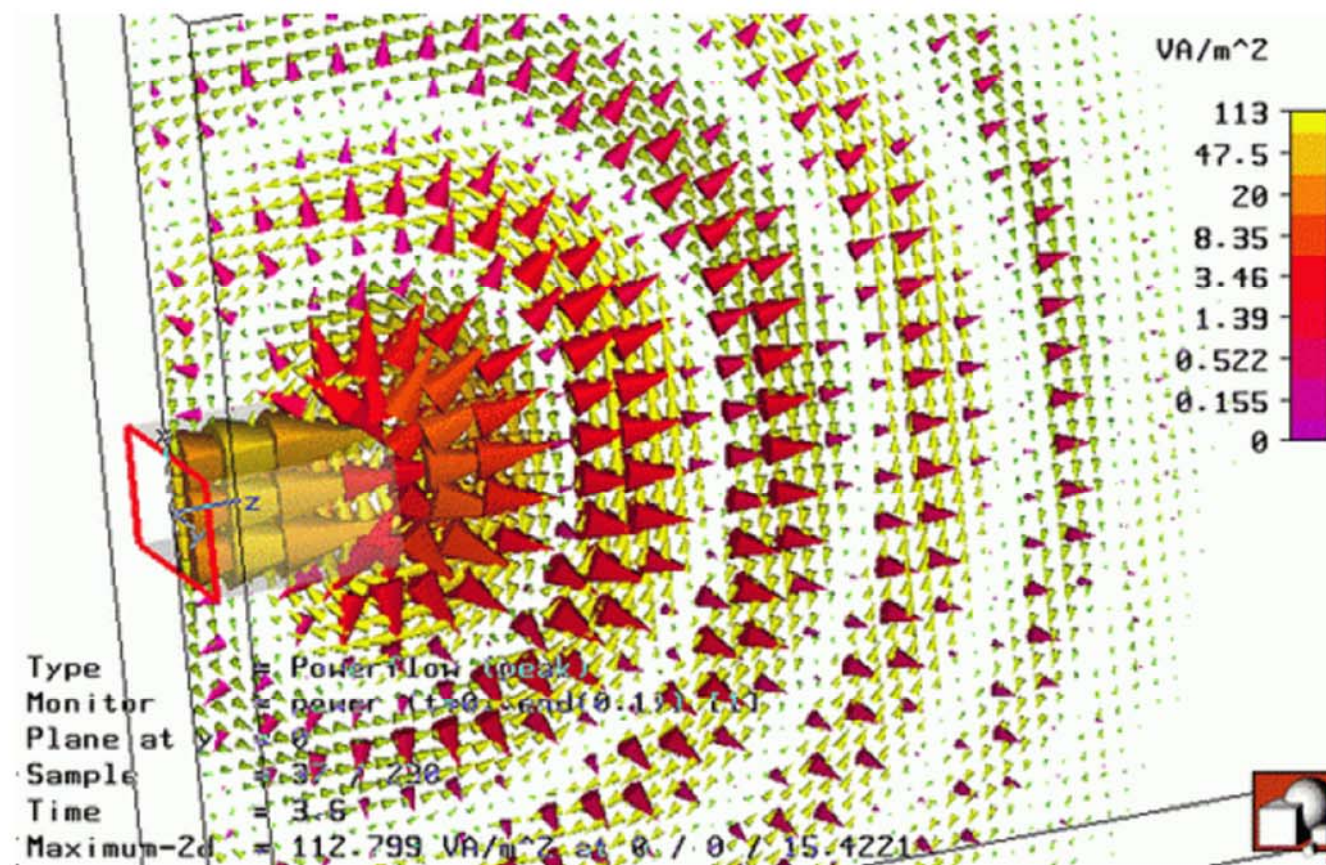


Radiation - H

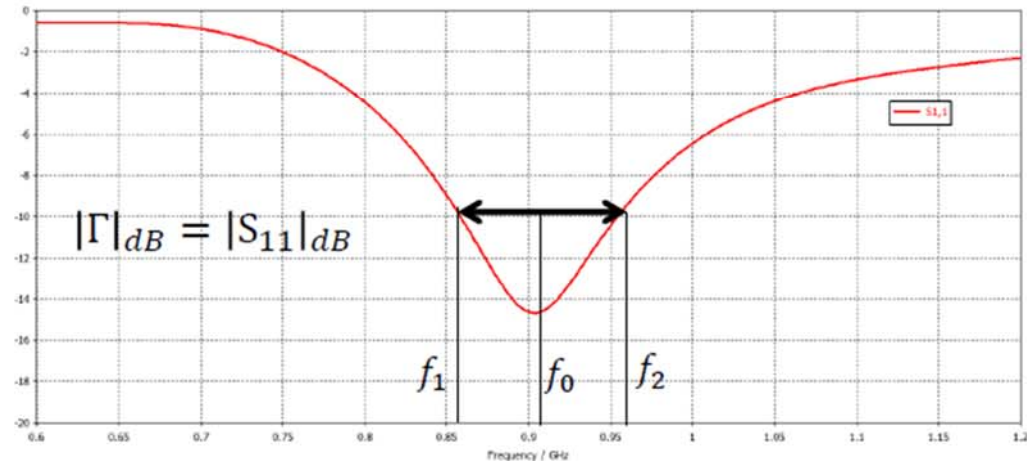
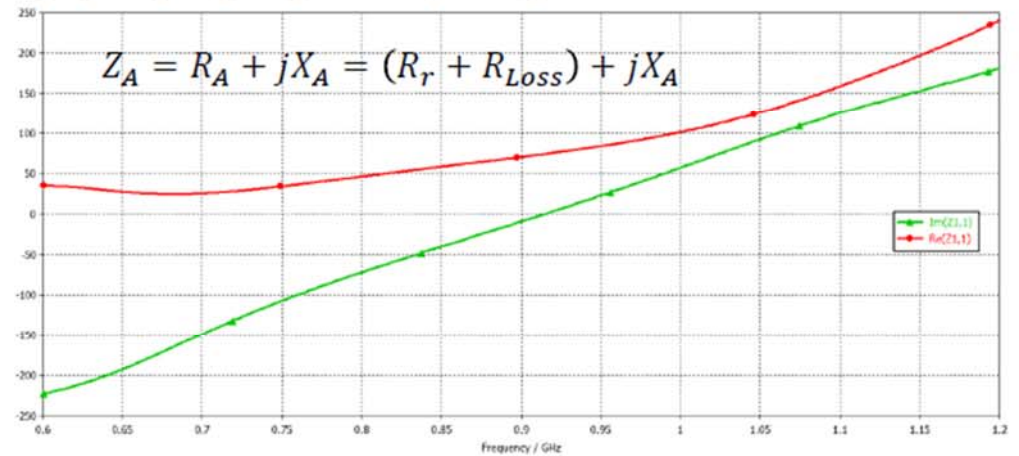
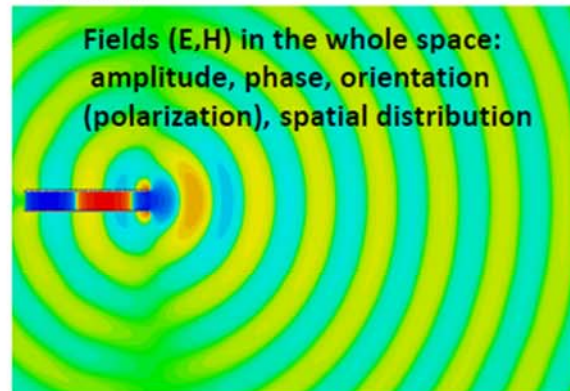
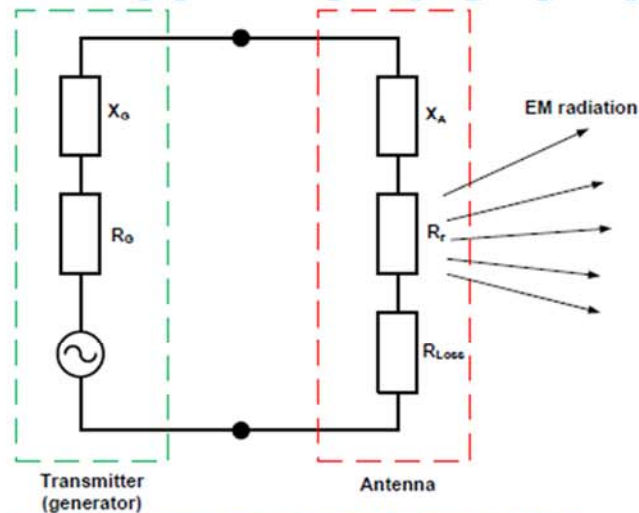


Radiation – power flow (Poynting vector)

$$\mathbf{S} = \mathbf{E} \times \mathbf{H}$$



Antenna as a circuit element



$$\Gamma = \frac{Z_A - Z_0}{Z_A + Z_0}$$

$$|\Gamma|_{dB} = 20 \log(|\Gamma|)$$

$$P_{reflected} = P_{input} |\Gamma|^2$$

$$VSWR = \frac{1 + |\Gamma|}{1 - |\Gamma|}$$

Relative impedance
bandwidth

$$BW = \frac{f_2 - f_1}{f_0} = \frac{f_2 - f_1}{\sqrt{f_1 f_2}}$$

Radiation efficiency

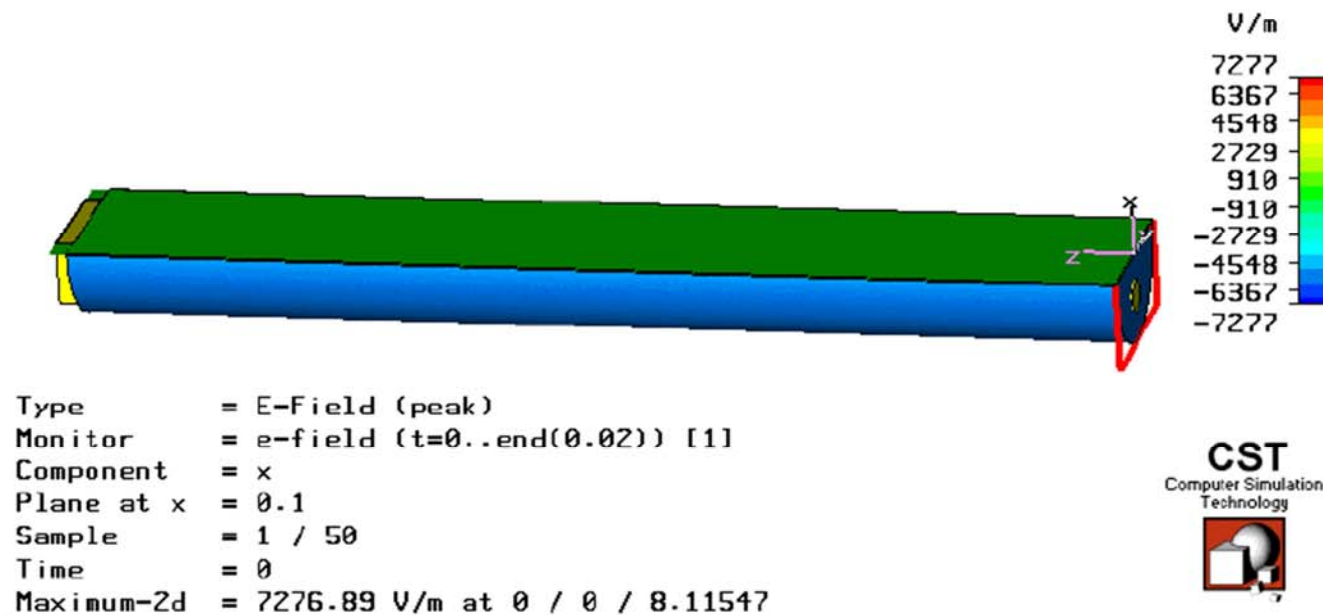
$$\eta_r = \frac{R_r}{R_r + R_{Loss}}$$

Matching

$$\Gamma = \frac{Z_A - Z_0}{Z_A + Z_0}$$

$$VSWR = \frac{1 + |\Gamma|}{1 - |\Gamma|}$$

$$|\Gamma|_{dB} = 20\log(|\Gamma|)$$



Elementary electric dipole

dz piece of electric current I is placed at origin in the z -axis, dipole length $L \ll \lambda$ ($L \rightarrow 0$)

$$\mathbf{A}(\mathbf{r}) = \frac{\mu}{4\pi} \int_V \mathbf{J}(\mathbf{r}') \frac{e^{-jkR}}{R} dV' \quad R = |\mathbf{r} - \mathbf{r}'| = \sqrt{x^2 + y^2 + (z - z')^2} = r \text{ for } L \rightarrow 0 \quad J = I\delta(x')\delta(y')\mathbf{z}_0$$

$$A_z(\mathbf{r}) = \frac{\mu I}{4\pi} \frac{e^{-jkr}}{r} \int_{-L/2}^{L/2} dz' = \frac{\mu IL}{4\pi} \frac{e^{-jkr}}{r}$$

Vector potential has the same orientation as current!
 $A_x = A_y = 0$

Transformation between rectangular and spherical vector components
 (it is convenient here to describe source points in cartesian and observation points in spherical coordinates respectively)

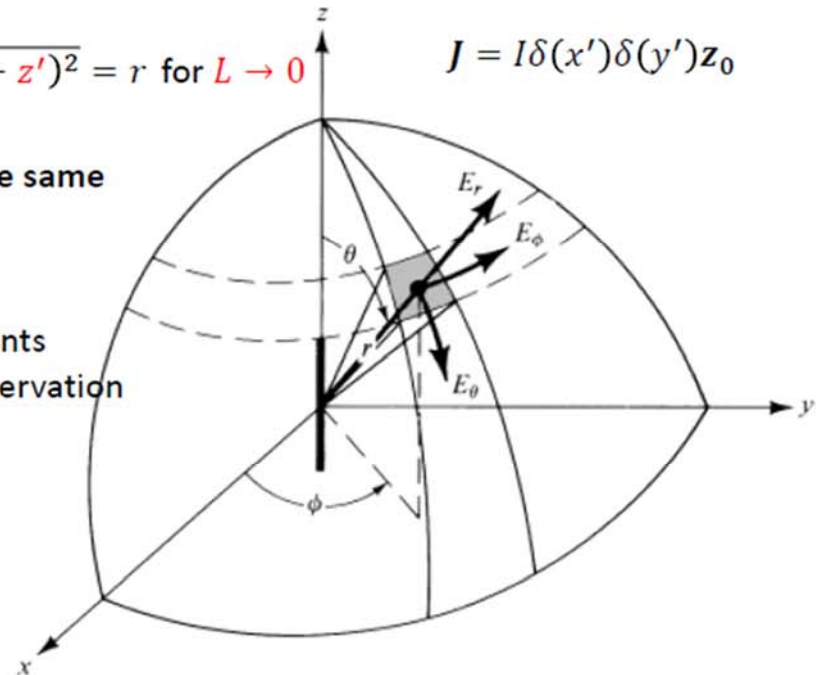
$$\begin{bmatrix} A_r \\ A_\theta \\ A_\phi \end{bmatrix} = \begin{bmatrix} \sin\theta \cos\phi & \sin\theta \sin\phi & \cos\theta \\ \cos\theta \cos\phi & \cos\theta \sin\phi & -\sin\theta \\ -\sin\phi & \cos\phi & 0 \end{bmatrix} \begin{bmatrix} A_x \\ A_y \\ A_z \end{bmatrix}$$

Spherical components of vector potentials in spherical coordinates:

$$A_r(\theta, \phi, r) = A_z \cos\theta = \frac{\mu IL}{4\pi} \frac{e^{-jkr}}{r} \cos\theta$$

$$A_\theta(\theta, \phi, r) = -A_z \sin\theta = -\frac{\mu IL}{4\pi} \frac{e^{-jkr}}{r} \sin\theta$$

$$A_\phi(\theta, \phi, r) = 0$$



$$\mathbf{E} = -j\omega\mathbf{A} - \frac{j}{\omega\mu\epsilon} \nabla(\nabla \cdot \mathbf{A})$$

$$\mathbf{H} = \frac{1}{\mu} \nabla \times \mathbf{A}$$

Elementary electric dipole

Components of electric and magnetic fields generated by an elementary electric dipole

$$H_{\phi}(r, \theta) = \frac{jkIL}{4\pi} \sin \theta \left[\frac{1}{r} + \frac{1}{jkr^2} \right] e^{-jkr}$$

$$E_r(r, \theta) = \frac{Z_0 IL}{2\pi} \cos \theta \left[\frac{1}{r^2} + \frac{1}{jkr^3} \right] e^{-jkr}$$

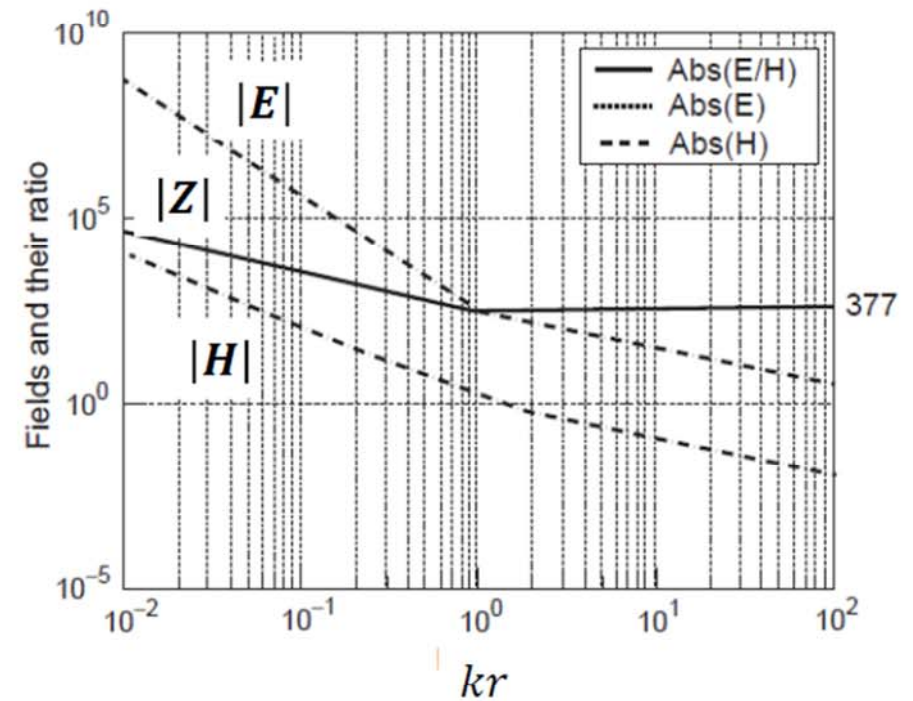
$$E_{\theta}(r, \theta) = \frac{jkZ_0 IL}{4\pi} \sin \theta \left[\frac{1}{r} + \frac{1}{jkr^2} - \frac{1}{k^2 r^3} \right] e^{-jkr}$$

$$H_r = H_{\theta} = 0$$

$$E_{\phi} = 0 \quad \text{Fields are not function of } \phi$$

General distance out of sources

These expressions are valid everywhere except on the source itself. In most cases we are interested only in the far-field, i.e. retain only $1/r$ terms. Generally, the radiated fields are vector functions of spherical coordinates r, θ, ϕ .

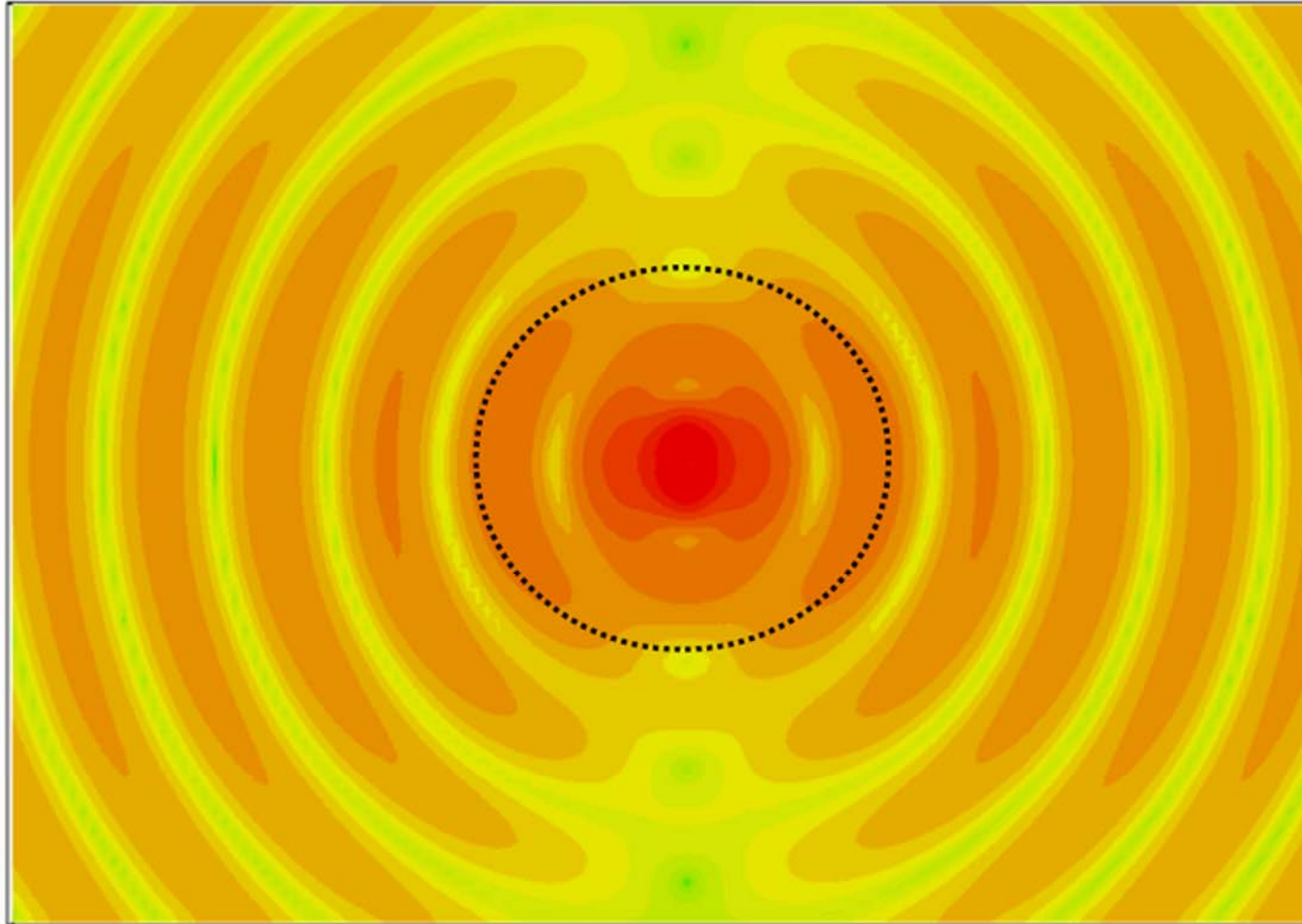


(distance)

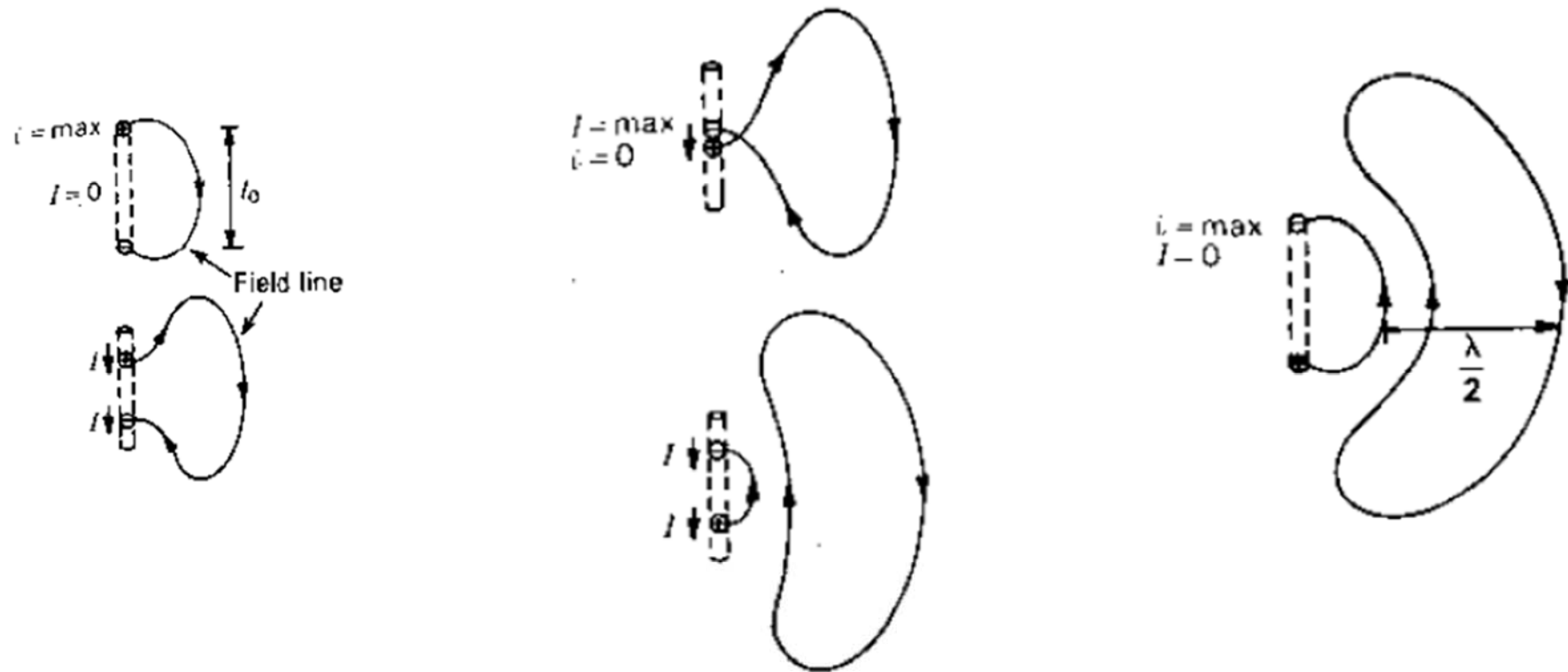
- E field is always greater than H field (its electric dipole)
- $kr = 1$ ($r = \lambda/2\pi$) is an important point (radian distance, radian sphere)
- $|Z| = \left| \frac{E}{H} \right| = 120\pi \cong 377\Omega$ if $kr > 1$

Elementary electric dipole

Short electric dipole ($L = \lambda/10$), magnitude of E-field intensity, radian-sphere $kr = 1$ ($r = \lambda/2\pi$) shown



Elementary electric dipole



Oscillating electric dipole consisting of two electric charges in simple harmonic motion, showing propagation of an electric field line and its detachment (radiation) from the dipole