

Lecture 7 Electrostatics (4)

1. Dielectric Polarization
2. Method of Images
3. Electrostatic Energy
4. Boundary Conditions
5. Coding Example

1. Dielectric Materials and Permittivity

■ Dielectric polarization

- Negative/positive charges oriented toward the positive/negative electrode
- Charges do not move freely in the material but reoriented.

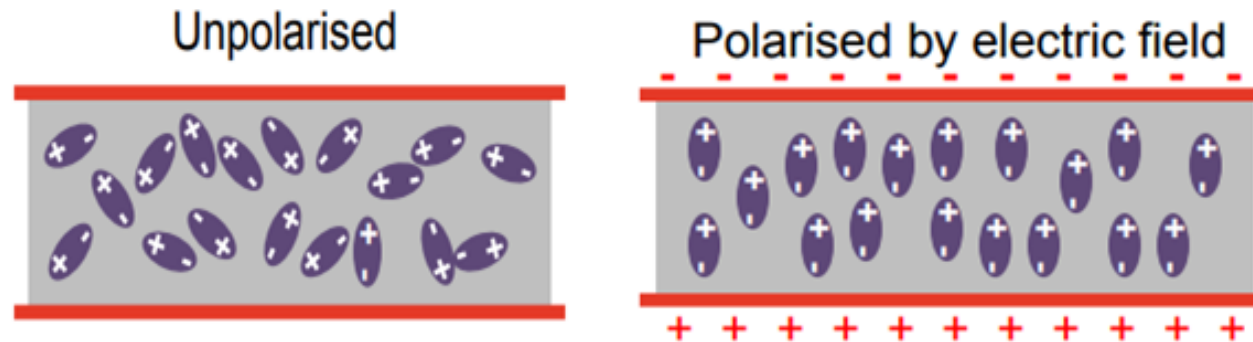
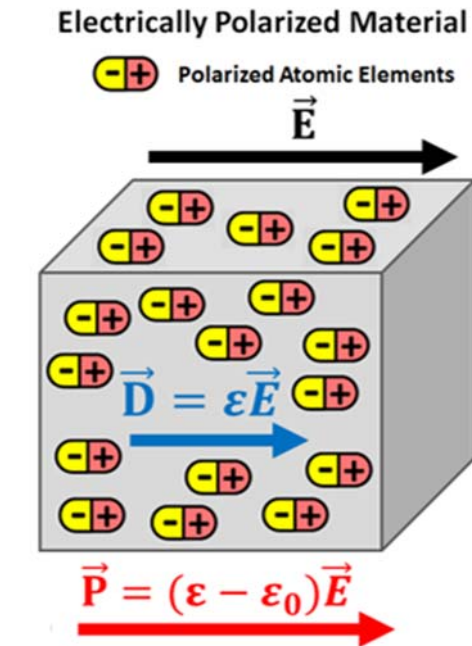


Fig. Polarization in materials [www.pcbdirectory.com]

■ Permittivity of dielectric materials

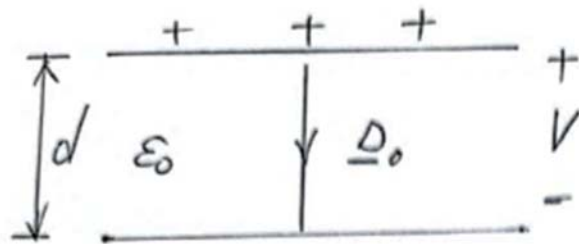


$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} = \epsilon_0 (1 + \chi_e) \mathbf{E} = \epsilon_r \epsilon_0 \mathbf{E} = \epsilon \mathbf{E}$$

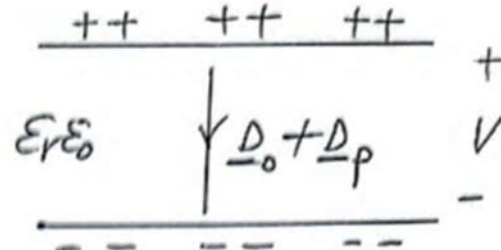
ϵ_0 : permittivity of vacuum, 8.854×10^{-12} F/m

$\epsilon = \epsilon_r \epsilon_0$: permittivity of dielectric materials

$\epsilon_r = \epsilon / \epsilon_0$: dielectric constant



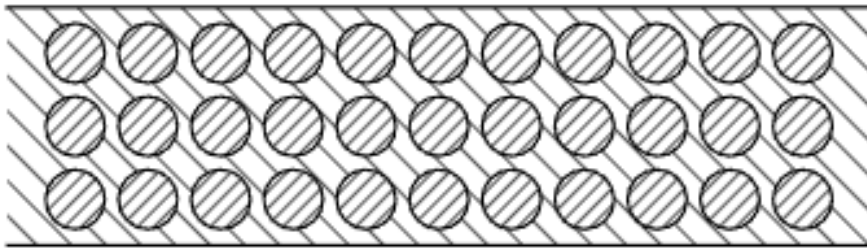
$$E = \frac{V}{d} \rightarrow D_0 = \epsilon_0 E = \epsilon_0 \frac{V}{d}$$



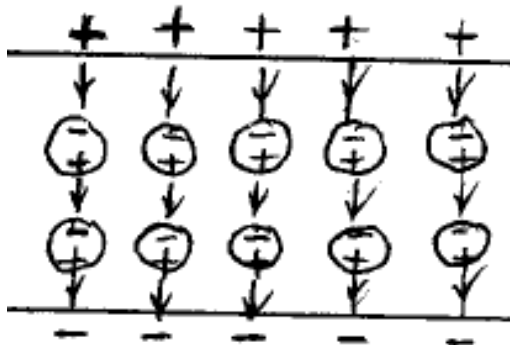
$$E = \frac{V}{d} \rightarrow D = D_0 + D_p = \epsilon_r \epsilon_0 E = \epsilon_r \epsilon_0 \frac{V}{d} > D_0$$

■ Dielectric polarization - Feynman's model

- A model of a dielectric: small conducting spheres embedded in an idealized insulator



- Charging of conducting spheres: top = negatively charged, bottom = positively charged



- Polarization (= charge separation) increases electric flux density in the material.

■ Polarization mechanisms in dielectric materials

- *Electron* displacements
- *Ion* displacements
- *Dipole* orientation
- *Space (interfacial) charge* displacements

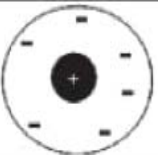
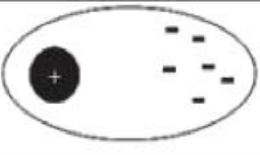
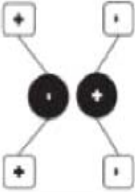
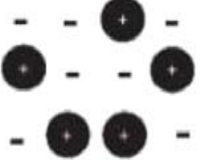
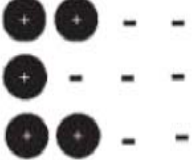
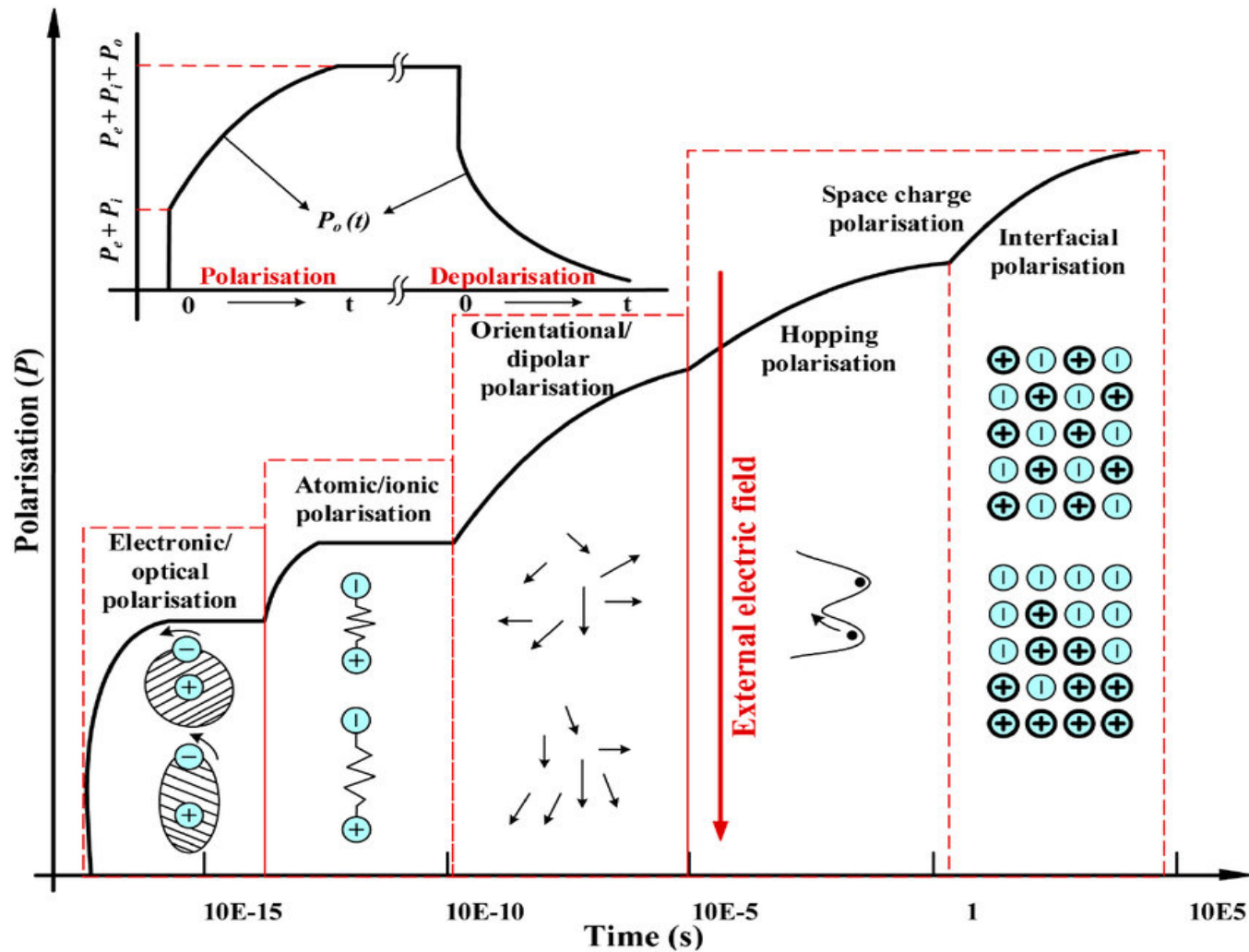
Polarization Mechanisms		
	No E field ($E = 0$)	← Local E field ← ($E \neq 0$)
Electronic		
Atomic or Ionic		
Orientation or Dipolar		
Interfacial		

Fig. Polarization mechanisms [Knowles]

■ Time constant of electric polarization

- Electric polarization: due to many different processes
- Polarization process: has a time constant (relaxation time)



■ Polarization charge density

- The effect of polarization is modelled using charge densities.
- *Polarization charges* cannot move around. → *Bound charges*

$\mathbf{P}$ = polarization $\equiv$ dipole moment per unit volume

$$V = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \hat{\mathbf{R}}}{R^2} \text{ (single dipole)}$$

$$V = \frac{1}{4\pi\epsilon_0} \int_V \frac{\mathbf{P} \cdot \hat{\mathbf{R}}}{R^2} dV' \text{ (distributed dipoles)}$$

$$V = \frac{1}{4\pi\epsilon_0} \oint_S \frac{\mathbf{P} \cdot d\mathbf{S}'}{R} + \frac{1}{4\pi\epsilon_0} \int_V \frac{(-\nabla' \cdot \mathbf{P})}{R} dV'$$

$\sigma_b \equiv \mathbf{P} \cdot \hat{\mathbf{n}}$ (polarized bound surface charge density)

$\rho_b \equiv -\nabla' \cdot \mathbf{P}$ (polarized bound volume charge density)

- Electric flux density in dielectric materials

$$\rho = \rho_b + \rho_f$$

$$\varepsilon_0 \nabla \cdot \mathbf{E} = \rho = \rho_b + \rho_f = -\nabla \cdot \mathbf{P} + \rho_f \quad (\text{Gauss's law})$$

$$\nabla \cdot (\varepsilon_0 \mathbf{E} + \mathbf{P}) = \rho_f$$

$$\mathbf{D} \equiv \varepsilon_0 \mathbf{E} + \mathbf{P} \quad (\text{electric flux density})$$

$$\nabla \cdot \mathbf{D} = \rho_f$$

■ Permittivity of dielectric materials

$$\mathbf{D} \equiv \varepsilon_0 \mathbf{E} + \mathbf{P} = \varepsilon_0 \mathbf{E} + \varepsilon_0 \chi_e \mathbf{E} = \varepsilon_0 (1 + \chi_e) \mathbf{E} = \varepsilon \mathbf{E}$$

$$\varepsilon \equiv \varepsilon_0 (1 + \chi_e) \text{ (permittivity)}$$

$$\varepsilon_r = 1 + \chi_e = \frac{\varepsilon}{\varepsilon_0} \text{ (dielectric constant or relative permittivity)}$$

Material	Dielectric Constant	Material	Dielectric Constant
Vacuum	1	Benzene	2.28
Helium	1.000065	Diamond	5.7-5.9
Neon	1.00013	Salt	5.9
Hydrogen (H ₂)	1.000254	Silicon	11.7
Argon	1.000517	Methanol	33.0
Air (dry)	1.000536	Water	80.1
Nitrogen (N ₂)	1.000548	Ice (-30° C)	104
Water vapor (100° C)	1.00589	KTaNbO ₃ (0° C)	34,000

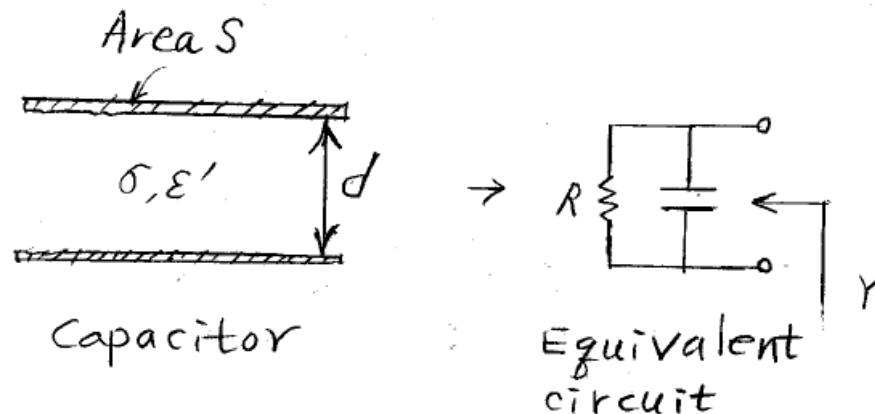
■ Dielectric loss

$$Y = G + j\omega C = \frac{\sigma S}{d} + j\omega \varepsilon \frac{S}{d} = j\omega \frac{S}{d} (\varepsilon' - j\varepsilon'') \rightarrow \varepsilon'' = \frac{\sigma}{\omega}$$

$$\omega = 2\pi f$$

$$G = \frac{\sigma S}{d}, C = \varepsilon' \frac{S}{d}$$

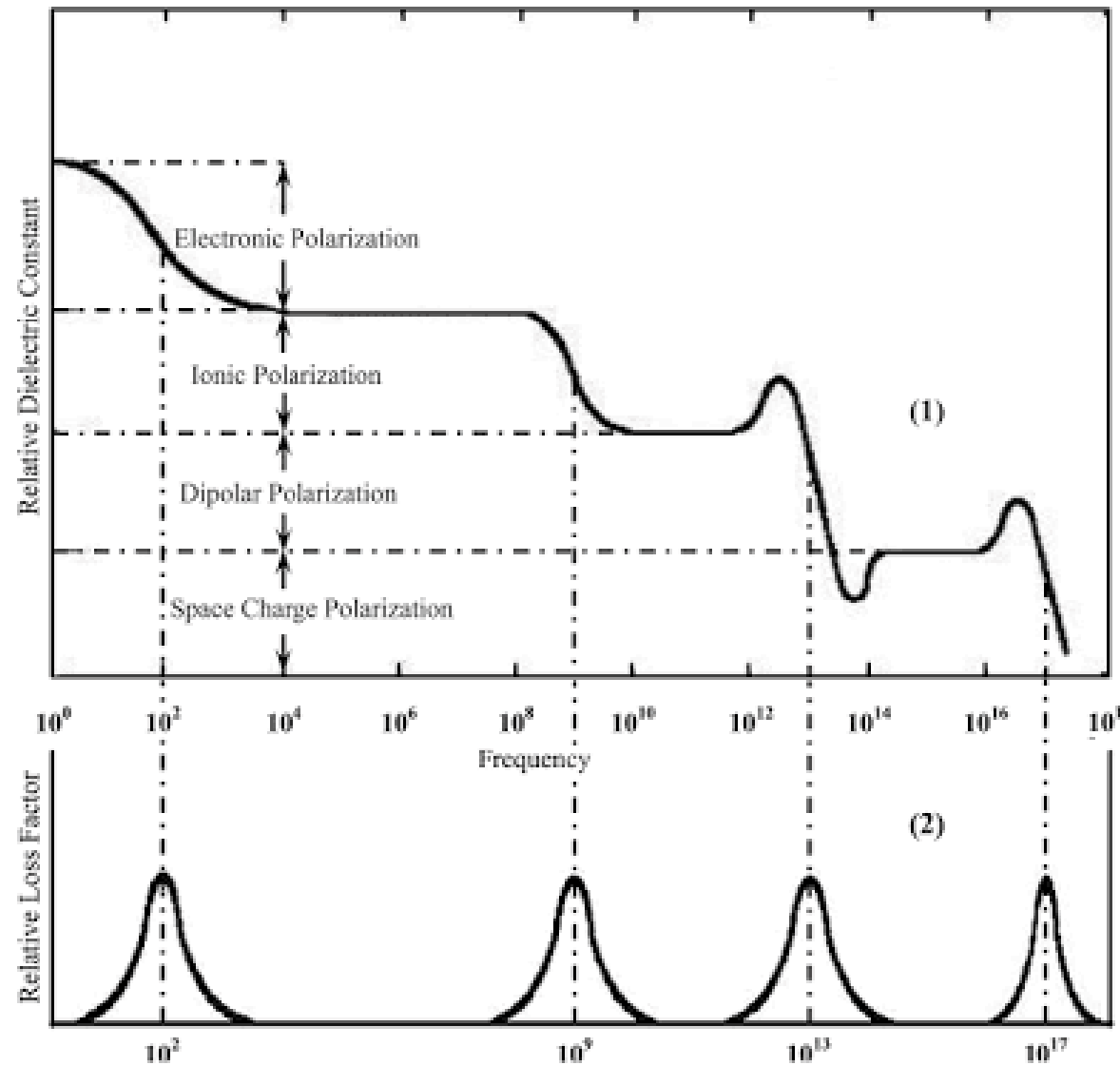
$$Q = CV, i = \frac{dQ}{dt} = C \frac{dV}{dt} \rightarrow I = j\omega CV \rightarrow Y = \frac{I}{V} = j\omega C$$



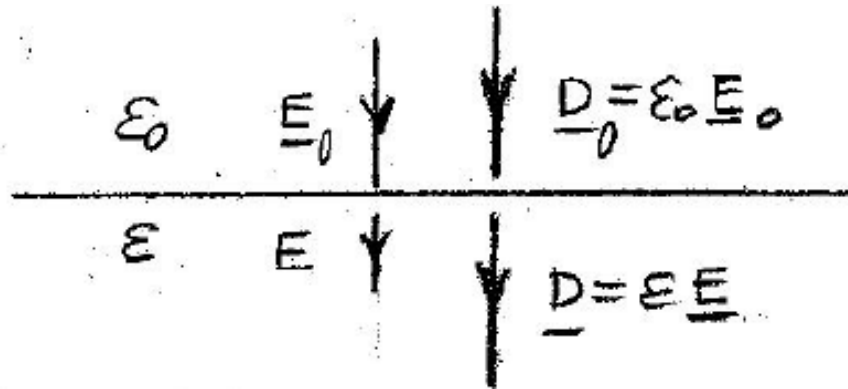
$$\varepsilon = \varepsilon' - j\varepsilon'' = \varepsilon_0 (\varepsilon'_r - j\varepsilon''_r) \text{ (complex permittivity)}$$

$$\tan \delta = \frac{\varepsilon''}{\varepsilon'} \text{ (loss tangent)}$$

- Permittivity dispersion: frequency response



- Example: dielectric in an external field



$\mathbf{D} = \mathbf{D}_0$ (continuity of electric flux; Gauss's law)

$$\mathbf{E} = \frac{\epsilon_0}{\epsilon} \mathbf{E}_0$$

$$|\mathbf{E}| < |\mathbf{E}_0|$$

2. Method of Images

$\mathbf{r} = (x, y, z)$: field point

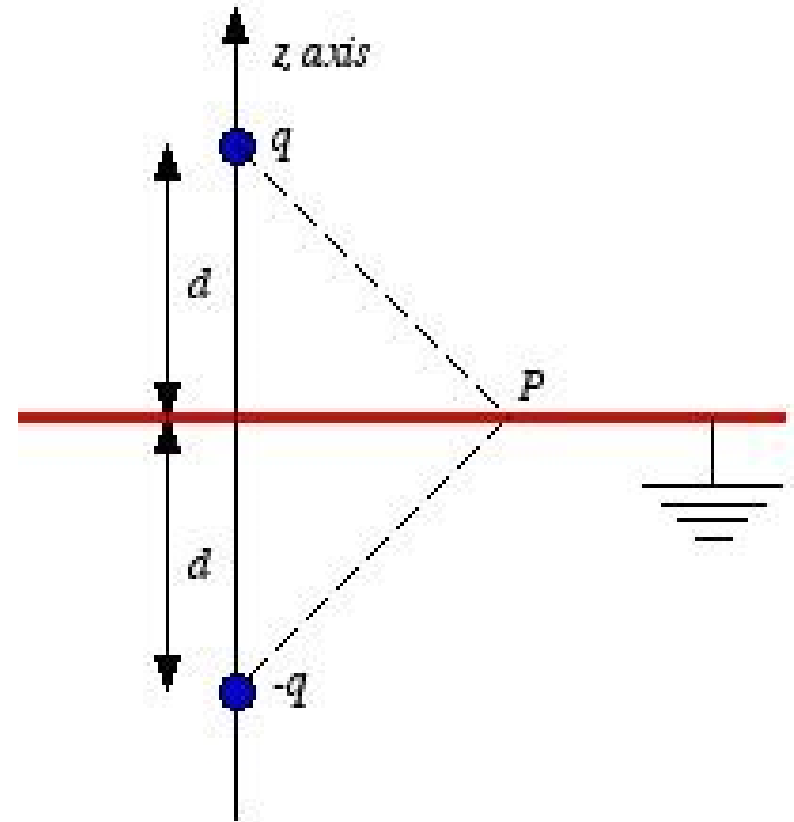
q at $\mathbf{r}_1 = (0, 0, d)$

$-q$ at $\mathbf{r}_2 = (0, 0, -d)$

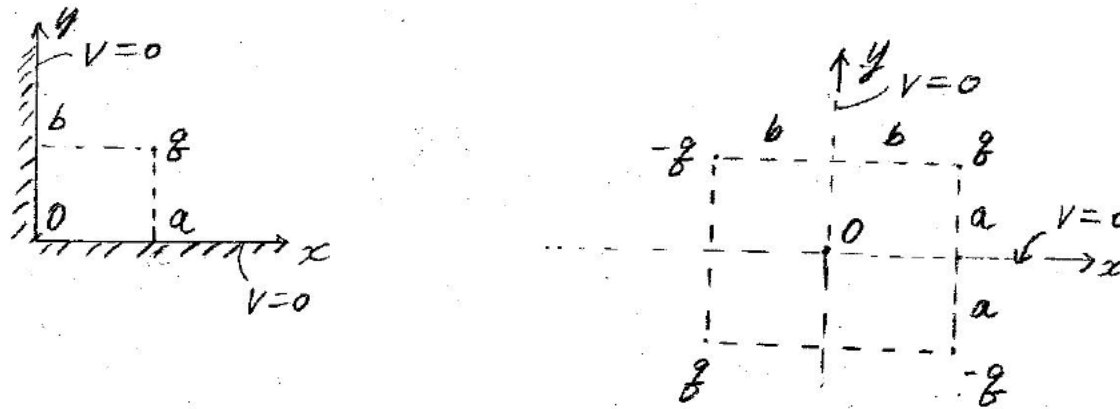
$\mathbf{R}_1 = \mathbf{r} - \mathbf{r}_1$, $\mathbf{R}_2 = \mathbf{r} - \mathbf{r}_2$

$$V(\mathbf{r}) = V_q + V_{-q} = \frac{1}{4\pi\epsilon} \frac{q}{R_1} - \frac{1}{4\pi\epsilon} \frac{q}{R_2}$$

$$\mathbf{E} = -\nabla V = \frac{1}{4\pi\epsilon} \frac{q\hat{\mathbf{R}}_1}{R_1^3} - \frac{1}{4\pi\epsilon} \frac{q\hat{\mathbf{R}}_2}{R_2^3}$$



- Example: a point source in a right-angle conductor



$\mathbf{r} = (x, y, z)$: field point

q at $\mathbf{r}_1 = (a, b, 0)$: the original point charge

$-q$ at $\mathbf{r}_2 = (-a, b, 0)$: an image charge

q at $\mathbf{r}_3 = (-a, -b, 0)$: an image charge

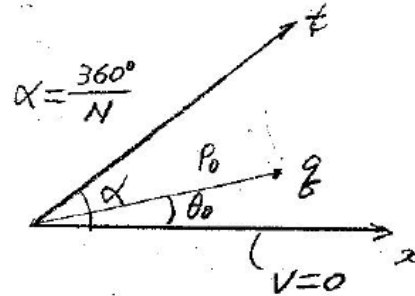
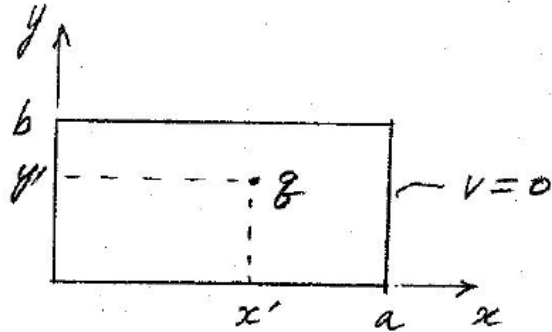
$-q$ at $\mathbf{r}_4 = (a, -b, 0)$: an image charge

$$\mathbf{R}_1 = \mathbf{r} - \mathbf{r}_1, \mathbf{R}_2 = \mathbf{r} - \mathbf{r}_2, \mathbf{R}_3 = \mathbf{r} - \mathbf{r}_3, \mathbf{R}_4 = \mathbf{r} - \mathbf{r}_4$$

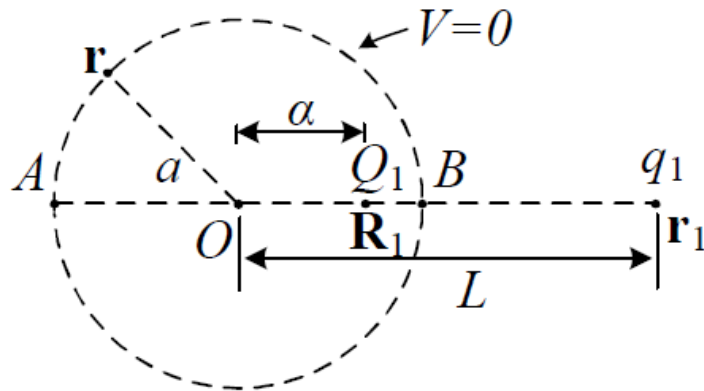
$$V(r) = \frac{q}{4\pi\epsilon} \left(\frac{1}{R_1} - \frac{1}{R_2} + \frac{1}{R_3} - \frac{1}{R_4} \right)$$

$$\mathbf{E} = -\nabla V = \frac{q}{4\pi\epsilon} \left(\frac{\hat{\mathbf{R}}_1}{R_1^3} - \frac{\hat{\mathbf{R}}_2}{R_2^3} + \frac{\hat{\mathbf{R}}_3}{R_3^3} - \frac{\hat{\mathbf{R}}_4}{R_4^3} \right)$$

- Example: problems requiring an infinite number of images



- Example: a point charge outside a conducting sphere



$$Q_1 = -\frac{a}{L} q_1$$

$$\alpha = \frac{a^2}{L}$$

3. Electrostatic Energy

■ Energy stored in N point charges

q_1
 $\cdot$
 r_1

q_1 only: $W = q_1 V_1 = 0$ since no field exists before

$$q_1, q_2 : W = W_2 = q_2 V_2 = q_2 \frac{q_1}{4\pi\epsilon |\mathbf{r}_2 - \mathbf{r}_1|}$$

q_1 q_2
 $\cdot$ $\cdot$
 r_1 r_2

$$q_1, q_2, q_3 : W = W_2 + W_3 = q_2 \frac{q_1}{4\pi\epsilon |\mathbf{r}_2 - \mathbf{r}_1|} + q_3 \left(\frac{q_1}{4\pi\epsilon |\mathbf{r}_3 - \mathbf{r}_1|} + \frac{q_2}{4\pi\epsilon |\mathbf{r}_3 - \mathbf{r}_2|} \right)$$

$$q_1, q_2, q_3, \dots, q_N : W = W_2 + W_3 + \dots + W_N = \frac{1}{4\pi\epsilon} \sum_{i=1}^N \left(q_i \sum_{j=1}^{i-1} \frac{q_j}{|\mathbf{r}_i - \mathbf{r}_j|} \right) = \frac{1}{2} \sum_{i=1}^N q_i V_i, \quad V_i = \frac{1}{4\pi\epsilon} \sum_{\substack{j=1 \\ j \neq i}}^N \frac{q_j}{|\mathbf{r}_i - \mathbf{r}_j|}$$

q_1 q_2 q_3
 $\cdot$ $\cdot$ $\cdot$
 r_1 r_2 r_3

q_1 q_2
 $\cdot$ $\cdot$
 r_1 r_2

q_N
 $\cdot$
 r_N

3. Electrostatic Energy

- Energy stored in distributed charges

$$W = \frac{1}{2} \sum_{i=1}^N q_i V_i$$

$$q_i = \rho_L dL = \rho_s dS = \rho_v dV$$

$$V_i = \phi(\mathbf{r})$$

$$\sum \rightarrow \int$$

$$W = \frac{1}{2} \int_L \rho_L \phi dL = \frac{1}{2} \int_S \rho_s \phi dS = \frac{1}{2} \int_V \rho_v \phi dV$$

- Electrostatic energy in terms of electric field

$$W = \int_V \frac{1}{2} \rho_v V dV = \int_V \frac{1}{2} (\nabla \cdot \mathbf{D}) V dV \quad (\because \rho_v = \nabla \cdot \mathbf{D})$$

$$\nabla \cdot (V \mathbf{D}) = V (\nabla \cdot \mathbf{D}) + \mathbf{D} \cdot (\nabla V) \rightarrow (\nabla \cdot \mathbf{D}) V = \nabla \cdot (V \mathbf{D}) - \mathbf{D} \cdot (\nabla V) = \nabla \cdot (V \mathbf{D}) + \mathbf{D} \cdot \mathbf{E}$$

$$W = \int_V \frac{1}{2} \nabla \cdot (V \mathbf{D}) dV + \int_V \frac{1}{2} (\mathbf{D} \cdot \mathbf{E}) dV = \oint_S \frac{1}{2} V \mathbf{D} \cdot d\mathbf{S} + \int_V \frac{1}{2} (\mathbf{D} \cdot \mathbf{E}) dV$$

$$w_s = \frac{1}{2} V \mathbf{D} \cdot \mathbf{n} \quad (\text{surface density of electric energy, J/m}^2)$$

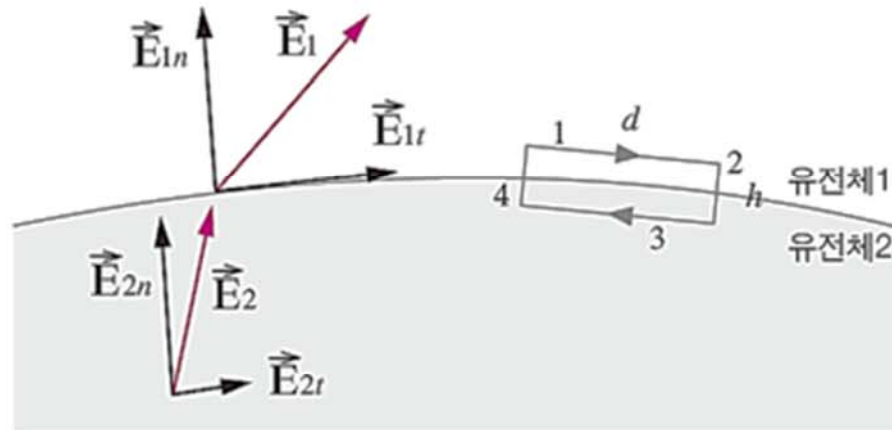
$$w_v = \frac{1}{2} \mathbf{D} \cdot \mathbf{E} \quad (\text{volume density of electric energy, J/m}^3)$$

$$\oint_S \frac{1}{2} V \mathbf{D} \cdot d\mathbf{S} = 0 \text{ as } V \rightarrow \infty$$

$$W = \int_V \frac{1}{2} (\mathbf{D} \cdot \mathbf{E}) dV$$

4. Boundary Conditions

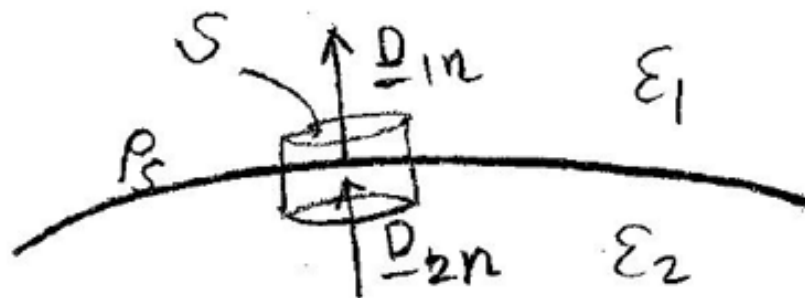
- Boundary conditions on a dielectric interface



$$\begin{aligned}\oint \vec{E} \cdot d\vec{l} &= \int_1 \vec{E}_{1t} \cdot d\vec{l} + \int_2 \vec{E}_n \cdot d\vec{l} + \int_3 \vec{E}_{2t} \cdot d\vec{l} + \int_4 \vec{E}_n \cdot d\vec{l} \\ &= E_{1t}d - (E_{1n} + E_{2n})\frac{h}{2} - E_{2t}d + (E_{1n} + E_{2n})\frac{h}{2} \\ &= (E_{1t} - E_{2t})d = 0\end{aligned}$$



$$E_{1t} = E_{2t}$$



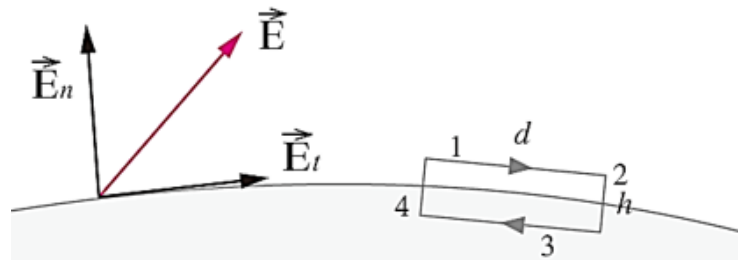
$$\oint \mathbf{D} \cdot d\mathbf{S} = (D_{1n} - D_{2n})S = Q = \rho_s S$$

$$D_{1n} - D_{2n} = \rho_s$$

$$\frac{E_{1n}}{\epsilon_1} - \frac{E_{2n}}{\epsilon_2} = \rho_s$$

$$\rho_s = 0 \rightarrow D_{1n} = D_{2n}, \quad \frac{E_{1n}}{\epsilon_1} = \frac{E_{2n}}{\epsilon_2}$$

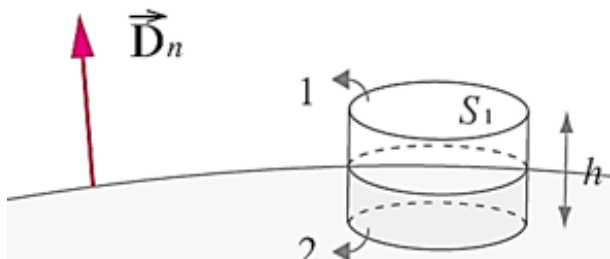
- Boundary conditions on a dielectric-conductor interface



$$\vec{E} = \vec{E}_t + \vec{E}_n$$

$$\begin{aligned} \oint \vec{E} \cdot d\vec{l} &= \int_1 \vec{E}_t \cdot d\vec{l} + \int_2 \vec{E}_n \cdot d\vec{l} + \int_3 \vec{E}_t \cdot d\vec{l} + \int_4 \vec{E}_n \cdot d\vec{l} \\ &= E_t d - E_n \frac{h}{2} + 0 + E_n \frac{h}{2} = E_t d \end{aligned}$$

$$E_t = 0$$

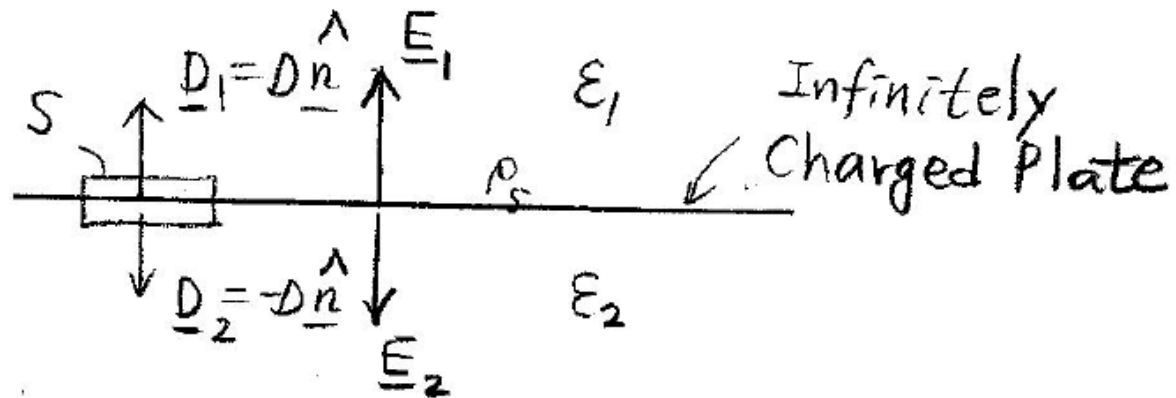


$$\begin{aligned} \oint \vec{D} \cdot d\vec{S} &= \int_{S_1} \vec{D}_n \cdot d\vec{S} + \int_{S_2} \vec{D}_n \cdot d\vec{S} \\ &= \epsilon_0 E_n S \\ &= Q \end{aligned}$$

$$E_n = \frac{Q}{\epsilon_0 S} = \frac{\rho_s}{\epsilon_0}$$

$$\vec{E} = \frac{\rho_s}{\epsilon_0} \hat{n} \text{ (V/m)}$$

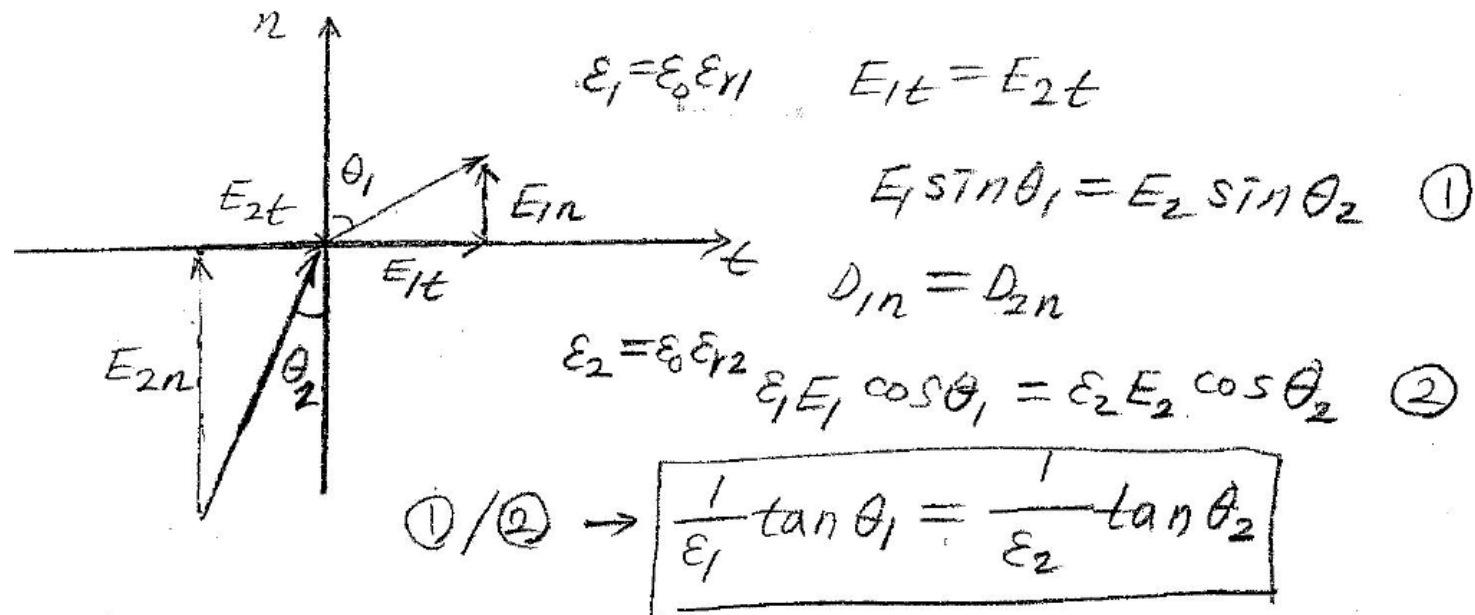
- Example: infinite charged plate



$$\oint \underline{D} \cdot d\underline{S} = Q \rightarrow 2DS = S\rho_s \rightarrow D = \frac{\rho_s}{2}$$

$$\underline{E} = \frac{\underline{D}}{\epsilon} \rightarrow \underline{E}_1 = \frac{\rho_s}{2\epsilon_1} \hat{n}, \quad \underline{E}_2 = -\frac{\rho_s}{2\epsilon_2} \hat{n}$$

- Example: angles of electric field lines on a dielectric interface



$$\epsilon_1 > \epsilon_2 \rightarrow \theta_1 > \theta_2$$

$$E_1 = E_2 \frac{\sin \theta_2}{\sin \theta_1} < E_2$$

$$\epsilon_{r1} = 80 \text{ (water)}, \epsilon_{r2} = 1 \text{ (air)}$$

$$\theta_2 = 15^\circ \rightarrow \tan \theta_1 = \frac{\epsilon_1}{\epsilon_2} \tan \theta_2 = 80 \tan \theta_2, \theta_1 = 87.3^\circ$$

$$E_2 = 1 \text{ (V/m)} \rightarrow E_1 = E_2 \frac{\sin \theta_2}{\sin \theta_1} = 0.26 \text{ (V/m)}$$

5. Coding Example

■ Example: angles of electric field lines on a dielectric interface

Write down a Python code for the stored energy for N point charges.

Input data:

N : number of point charges

q_i (C), (x_i, y_i, z_i) m : i -th point charge and position

Output data:

W (J): stored energy

Lecture 7 - Python coding example: Stored energy in n point charges

Input data:

n = number of point charges (maximum 100)

q_i = charge (C) of the i -th charge

(x_i, y_i, z_i) = position (m) of the i -th charge

Output data:

w = stored energy (J)

```
from math import *
```

```
e0=8.853e-12;pi=3.141593
```

```
x=[0]*100;y=[0]*100;z=[0]*100;q=[0]*100
```

```
while True:
```

```
    n=int(input('Number of point charges='))
```

```
    for i in range(n):
```

```
        print('For charge no.',i)
```

```
        q[i]=float(input('Charge: q(C)='))
```

```
        x[i],y[i],z[i]=map(float,input('Position: x,y,z(m)=').split())
```

```
    w=0
```

```
    for i in range(n):
```

```
        sum=0
```

```
        for j in range(n):
```

```
            if(j == i):
```

```
                continue
```

```
            else:
```

```
                sum=sum+q[j]/sqrt((x[i]-x[j])**2+(y[i]-y[j])**2+(z[i]-z[j])**2)
```

```
        w=w+q[i]*sum
```

```
w=w/(4*pi*e0)/2
```

```
print('Stored: W(J)=',w)
```

'''

Number of point charges=

3

For charge no. 0

Charge: $q(C)=$

$1e-10$

Position: $x,y,z(m)=$

1 1 3

For charge no. 1

Charge: $q(C)=$

$2e-4$

Position: $x,y,z(m)=$

3 -1 1

For charge no. 2

Charge: $q(C)=$

$4e-10$

Position: $x,y,z(m)=$

4 1 -4

Stored: $W(J)= 0.00018318585243$

Number of point charges=

'''

Fin
(End)